

# Rotational Magneto-Acousto-Electric Tomography: Theory and Experiments

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# Traditional modalities (please, don't take this slide literally!)

Type	Costs	Purpose	Mathematics	Instability
X-Ray	expensive	bone structure	Radon transform	mild
Gamma	expensive	blood flow	Attenuated Radon	stronger
Acoustic	cheap	soft tissue	Born/Rytov	mild
Microwave	cheap	high contrast	non-linear ?	???
MRI	very expensive		Fourier transform	mild
Impedance	very cheap	lung motion ?	divergence eq-n	very strong
Optical	cheap	small objects	diffusion eq-n ?	strong ?

# Hybrid methods: motivation

Conductivity in tumors is much higher than that in healthy tissues  
⇒ EM waves or currents yield high contrast.

**Electrical impedance tomography**, optical and microwave tomographies lead to strongly non-linear and ill-posed inverse problems = **BAD!**

Acoustic waves yield high resolution but the contrast is low.

**Idea:** Use **hybrid** techniques, couple ultrasound with EM field:

Thermo-Acoustic and Photo-Acoustic Tomography (TAT/PAT)

Ultrasound Modulated Optical Tomography (UMOT)

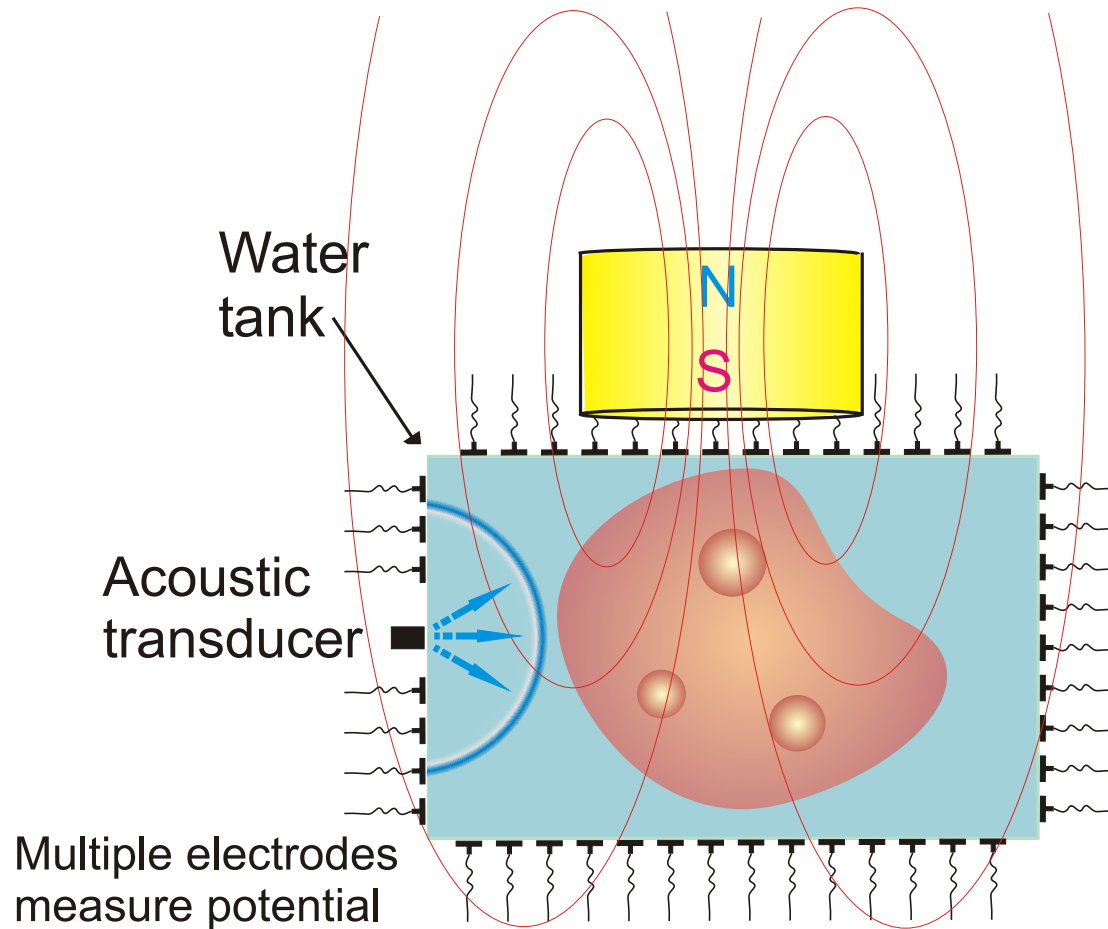
**Acousto-Electric Tomography (AET)**

Magneto-Acousto-Electric Tomography (MAET)

Magneto-Acoustic Tomography with Magnetic Induction (MAT-MI)

Some of these techniques are "theoretical"

# Lorentz Force Tomography (a.k.a MAET)



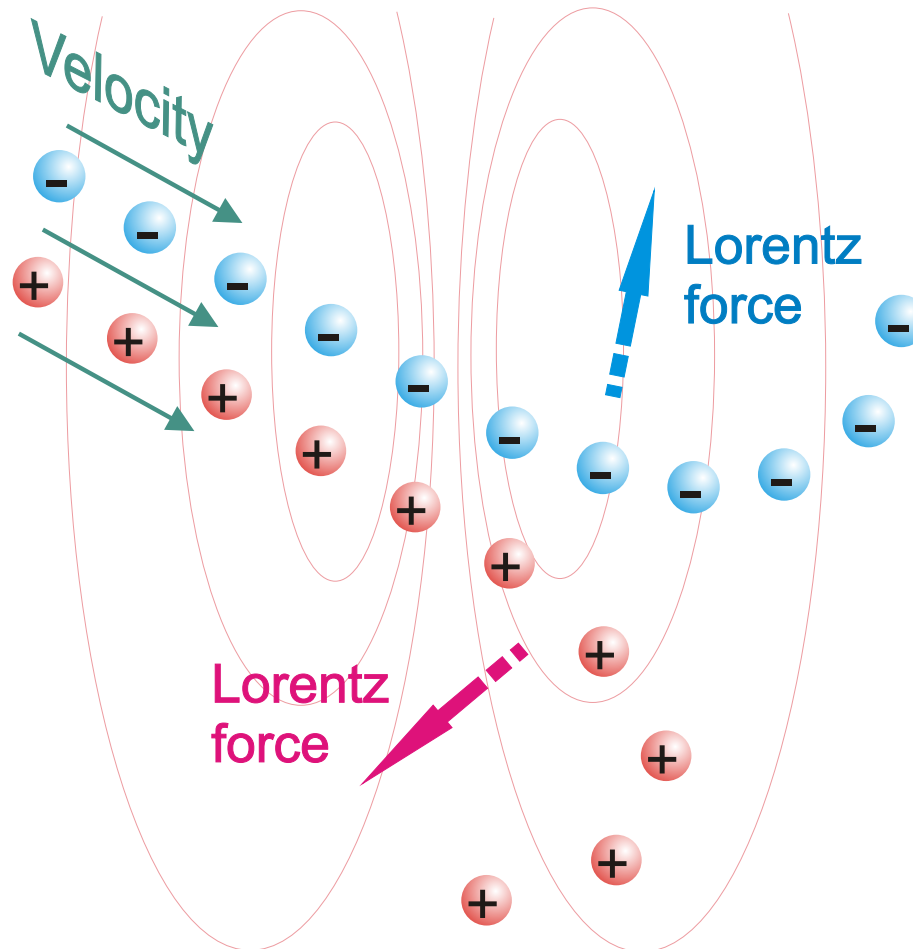
Ultrasound makes electrons and ions vibrate.

As a result, moving electrons and ions are separated by the Lorentz force.

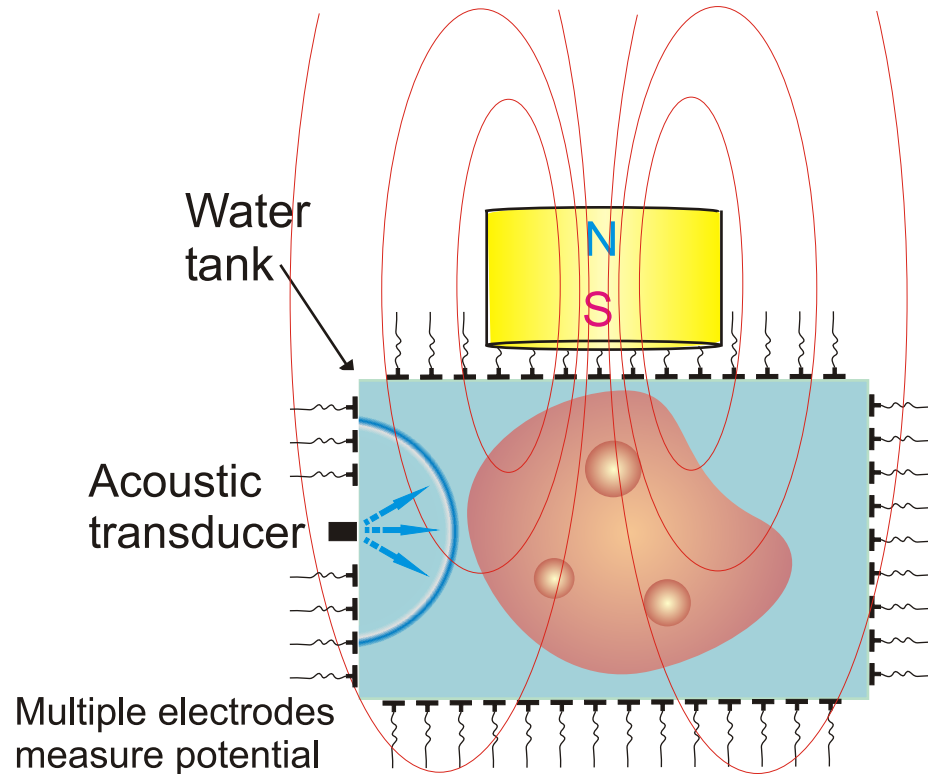
# What's the Lorentz Force?

In a magnetic field the Lorentz force pushes moving charges sideways

Positive and negative particles are pushed in the opposite directions



# MAET (a.k.a Lorentz Force Tomography)



Separated charges create an electric field that's picked up by the electrodes

With some clever mathematics one can reconstruct an image

# Previous work on MAET

The mathematics of MAET (partially explained below) is very promising

MAET signal has been demonstrated only in one-directional measurements

No truly tomographic MAET images have been obtained before

Our goal: to demonstrate the feasibility of a full-scale MAET

# Physics & mathematics of MAET

Tissue moving with velocity  $V(x, t)$  produces Lorentz currents  $J_L(x, t)$ :

$$J_L(x, t) = \sigma(x)B \times V(x, t)$$

There will also be Ohmic currents satisfying Ohm's law

$$J_O(x, t) = \sigma(x)\nabla u(x, t).$$

There are no sinks or sources, the total current is divergence-free

$$\nabla \cdot (J_L + J_O) = 0.$$

Thus

$$\nabla \cdot \sigma \nabla u = -\nabla \cdot (\sigma B \times V).$$

BC: the normal component of the total current  $J_L(x, t) + J_O(x, t)$  vanishes:

$$\left. \frac{\partial}{\partial n} u(z) \right|_{\partial\Omega} = -(B \times V(z)) \cdot n(z)$$



# Measuring functionals

At any given time  $t$  we measure potential  $u(z, t)$  at all  $z \in \partial\Omega$ .

Integrate boundary values of  $u$  with weight  $I(z)$  and get a functional  $M(t)$ :

$$M(t) = \int_{\partial\Omega} I(z)u(z, t)dA(z),$$

**Introduce lead currents = virtual currents**

Consider lead potential  $w_I(x)$  and lead current  $J_I(x) = \sigma(x)\nabla w_I(x)$ :

$$\begin{aligned}\nabla \cdot \sigma \nabla w_I(x) &= 0, \\ \left. \frac{\partial}{\partial n} w_I(z) \right|_{\partial\Omega} &= I(z).\end{aligned}$$

Then, using the second Green's identity (= **reciprocity** principle):

$$M(t) = \int_{\Omega} B \cdot J_I(x) \times V(x, t)dx$$

# Analyzing the velocity field

Assume that speed of sound  $c$  and density  $\rho$  are constant.

Then, velocity is the gradient of the velocity potential  $\varphi(x, t)$ :

$$V(x, t) = \frac{1}{\rho} \nabla \varphi(x, t),$$

where velocity potential  $\varphi(x, t)$  is the time anti-derivative of pressure  $p(x, t)$ :

$$p(x, t) = \frac{\partial}{\partial t} \varphi(x, t).$$

Substitute into equation for  $M(t)$  and integrate by parts:

$$M(t) = \frac{1}{\rho} B \cdot \left[ \int_{\partial\Omega} \varphi(z, t) J_I(z) \times n(z) dA(z) + \int_{\Omega} \varphi(x, t) \nabla \times J_I(x) dx \right]$$

Volumetric part shows that we measure components of **curl**  $J_I(x)$ !

$$\mathbf{curl} J_I(x) = \nabla \times [\sigma(x) \nabla w_I(x)] = \nabla \sigma(x) \times \nabla w_I(x) = \nabla \ln \sigma(x) \times J_I(x)$$

Notice: in the regions where  $\sigma(x)$  is constant,  $\mathbf{curl} J_I(x) = 0$ . No signal!

# Reconstruction procedure

If  $\varphi(x, t)$  could be focused into a point, i.e.  $\varphi(x, 0) = \delta(x - x_0)$ , then

$$M_{x_0}(0) = \frac{1}{\rho} B \cdot \left[ \int_{\Omega} \delta(x - x_0) \text{curl} J_I(x) dx \right] = \frac{1}{\rho} B \cdot \text{curl} J_I(x_0).$$

If three different directions of  $B$  are used, we have  $C(x_0) = \text{curl} J_I(x_0)$ !

## Chain of equations to solve:

Curl  $C$   $\rightarrow$  Current  $I$   $\rightarrow$   $\nabla \ln \sigma(x)$   $\rightarrow$  Conductivity  $\sigma(x)$ .

The second step comes from:

$$\nabla \ln \sigma \times J = C.$$

If we have two currents  $J^{(j)}(x)$ ,  $j = 1, 2$ , then solve for  $\nabla \ln \sigma$  at each  $x$

$$\begin{cases} \nabla \ln \sigma(x) \times J^{(1)}(x) = C^{(1)}(x) \\ \nabla \ln \sigma(x) \times J^{(2)}(x) = C^{(2)}(x) \end{cases} \cdot$$

# 3D MAET with ideal measurements

## To summarize:

The inverse problems for MAET with ideal measurements is

stable

(almost) explicitly solvable...

... by a linear algorithm

(see [Kunyansky, 2012])

This is very rare, and very promising from the engineering standpoint.

No experimental work for 3D MAET has been ever done.

# Something simpler: 2D MAET

Full 3-D scanner for MAET is difficult to build

We want to demonstrate the feasibility of MAET in a 2D setting

## Assumptions and approximations:

Everything is constant in the vertical direction ( $\vec{e}_z$ ).

Magnetic induction  $B = b\vec{e}_z$  (vertical and constant).

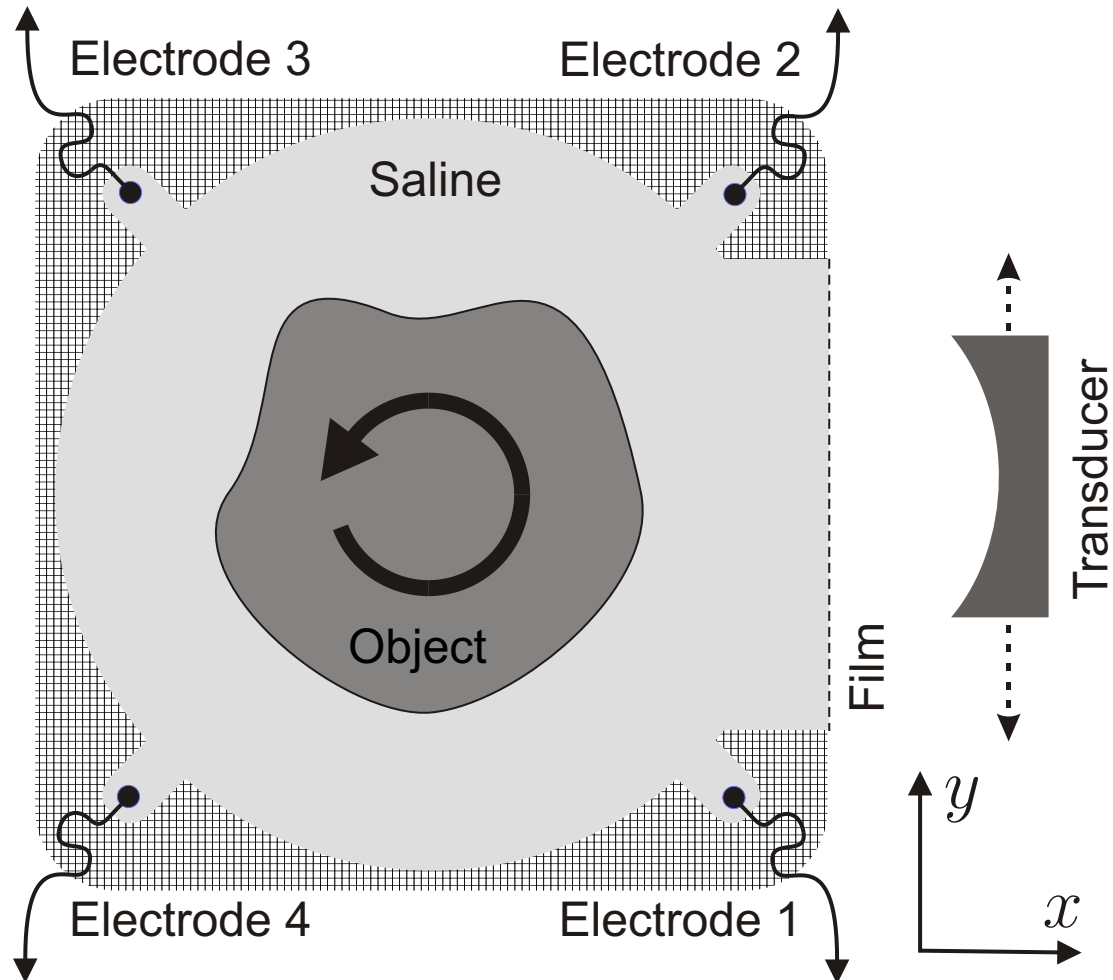
All the objects have vertical boundaries (generalized cylinders)

Electrodes are vertical lines

Then, all curls are vertical and parallel to  $B$  and we measure  $\frac{b}{\rho}\text{curl}_z J$

# 2D MAET scanner

A simple 2D MAET scanner, top view:



# Synthetic flat transducer

## Problem:

The exact time-dependent velocity field of a focusing transducer is complicated and difficult to measure.

Instead, we average all measurements corresponding to a fixed angular position of the object and varying vertical position of the transducer.

Due to linearity of the problem, this is equivalent to using a large, flat, vertically oriented sound-emitting source.

Approximately:

$$\varphi(t, x) = c_{\text{tr}} \delta(-x_{\text{tr}} + x_1 + ct),$$

and (if  $C(x)$  is the curl)

$$M(t) = \frac{bc_{\text{tr}}}{\rho} \int_{\Omega} \delta(-x_{\text{tr}} + x_1 + ct) C(x) dx = \frac{bc_{\text{tr}}}{\rho} \int_{\mathbb{R}} C [(x_{\text{tr}} - ct)\vec{e}_1 + s\vec{e}_2] ds,$$

Thus, we measure **Radon projections** of  $C(x)$  (integrals over vertical lines).

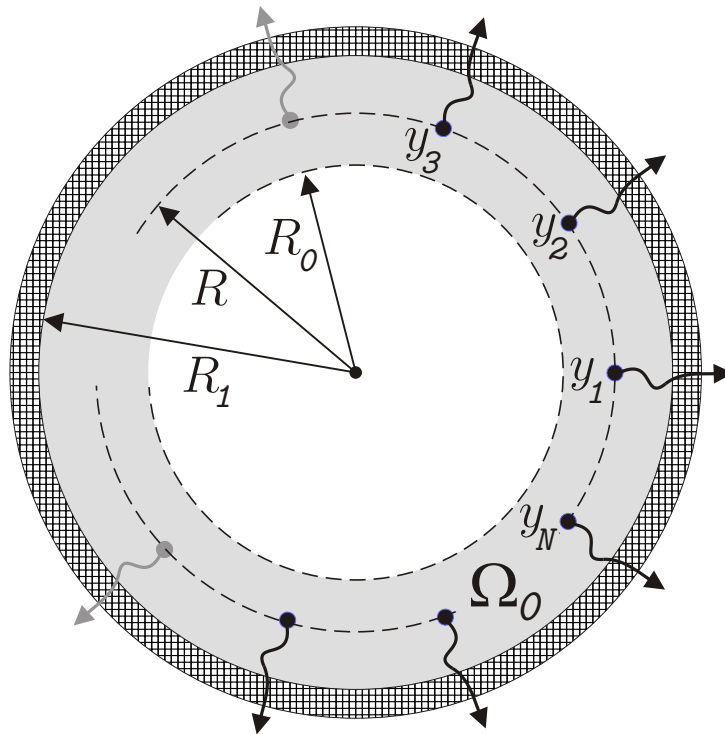
# MAET with a rotating object

## Problem:

We want to reconstruct curl  $C(x)$  from its Radon projections.

However, if the object rotates, but the electrodes are stationary, the currents change and the curl  $C(x)$  changes.

To resolve this, consider a round chamber, with  $N$  electrodes equispaced on a circle of radius  $R$ :





# Lead currents and potentials; round chamber

The lead potential corresponding to a set of weights  $\mathbf{W}$  equals:

$$w_{\mathbf{W}}(x) = w_{\mathbf{W}}^{\text{sing}}(x) + w_{\mathbf{W}}^{\text{smooth}}(x),$$

where  $w_{\mathbf{W}}^{\text{sing}}$  is defined as

$$w_{\mathbf{W}}^{\text{sing}}(x) \equiv h(x) + \frac{1}{2\pi\sigma_0} \sum_{j=1}^N W_j \ln |x - y_j|,$$

with  $h(x)$  harmonic in  $\Omega$  and such that

$$\frac{\partial}{\partial n} w_{\mathbf{W}}^{\text{sing}}(z) = 0, \quad z \in \partial\Omega.$$

Then  $w_{\mathbf{W}}^{\text{smooth}}(x)$  is the solution of the following BV problem:

$$\nabla \cdot \sigma \nabla w_{\mathbf{W}}^{\text{smooth}}(x) = -\chi(x) \nabla \cdot \sigma(x) \nabla w_{\mathbf{W}}^{\text{sing}}(x), \quad x \in \Omega.$$

$$\frac{\partial}{\partial n} w_{\mathbf{W}}^{\text{smooth}}(z) = 0, \quad z \in \partial\Omega,$$

Is  $w_{\mathbf{W}}^{\text{smooth}}$  a "scattering of incoming potential  $w_{\mathbf{W}}^{\text{sing}}$  by  $\sigma(x)$ " ?

# Synthetic lead currents

For an arbitrary unit vector  $\gamma = (\cos \alpha, \sin \alpha)$ ,  
define the set of weights  $\mathbf{W}^\gamma = (W_1^\gamma, \dots, W_N^\gamma)$  by the formula:

$$W_j^\gamma \equiv \frac{1}{N} \cos \left( \frac{2\pi(j-1)}{N} - \alpha \right).$$

Then

$$w_{\mathbf{W}^\gamma}^{\text{sing}}(x) \approx \beta \gamma \cdot x,$$

where  $\beta$  is a known constant.

This approximation converges exponentially in the limit  $N \rightarrow \infty$ .

Now, the corresponding lead potential  $w_{\mathbf{W}^\gamma}$  can be synthetically rotated, by simultaneously turning the object and adjusting angle  $\alpha$ .

**The rest of the problem is solved explicitly, as before**

# The case of a realistic piezoelectric transducer

## A big problem:

Widely used piezoelectric transducers do not reproduce low frequencies.

As a result, only a high-frequency component of  $C(x)$  can be reconstructed.

Lead currents cannot be reconstructed at all.

A **very crude** solution: use near-constant approximation of  $\sigma(x)$ . Then:

$$w_{\mathbf{W}\gamma}(x) \approx w_{\mathbf{W}\gamma}^{\text{sing}}(x) \approx \beta\gamma \cdot x,$$

$$J(x) \approx \sigma_0\beta\gamma,$$

$$C(x) = \sigma_0\beta\gamma^\perp \cdot \nabla \ln \sigma(x)$$

## Finally...

We use two orthogonal lead currents  $J^{(1)} \approx \sigma_0 \beta \gamma^{(1)}$  and  $J^{(2)} \approx \sigma_0 \beta \gamma^{(2)}$ , with curls  $C^{(1)}(x)$  and  $C^{(2)}(x)$ .

Then

$$C^{(1)}(x) \approx -\sigma_0 \beta \gamma^{(1)} \cdot \nabla \ln \sigma(x) = -\sigma_0 \beta \frac{\partial \ln \sigma(x)}{\partial \gamma^{(1)}},$$

$$C^{(2)}(x) \approx \sigma_0 \beta \gamma^{(2)} \cdot \nabla \ln \sigma(x) = \sigma_0 \beta \frac{\partial \ln \sigma(x)}{\partial \gamma^{(2)}},$$

and

$$\Delta \ln \sigma(x) \approx \frac{1}{\sigma_0 \beta} \left( \frac{\partial}{\partial \gamma^{(2)}} C^{(2)}(x) - \frac{\partial}{\partial \gamma^{(1)}} C^{(1)}(x) \right).$$

While this simple technique is based on a small perturbations of constant conductivity, in practice, it captures boundaries of material interfaces even if  $\sigma(x)$  is strongly nonuniform

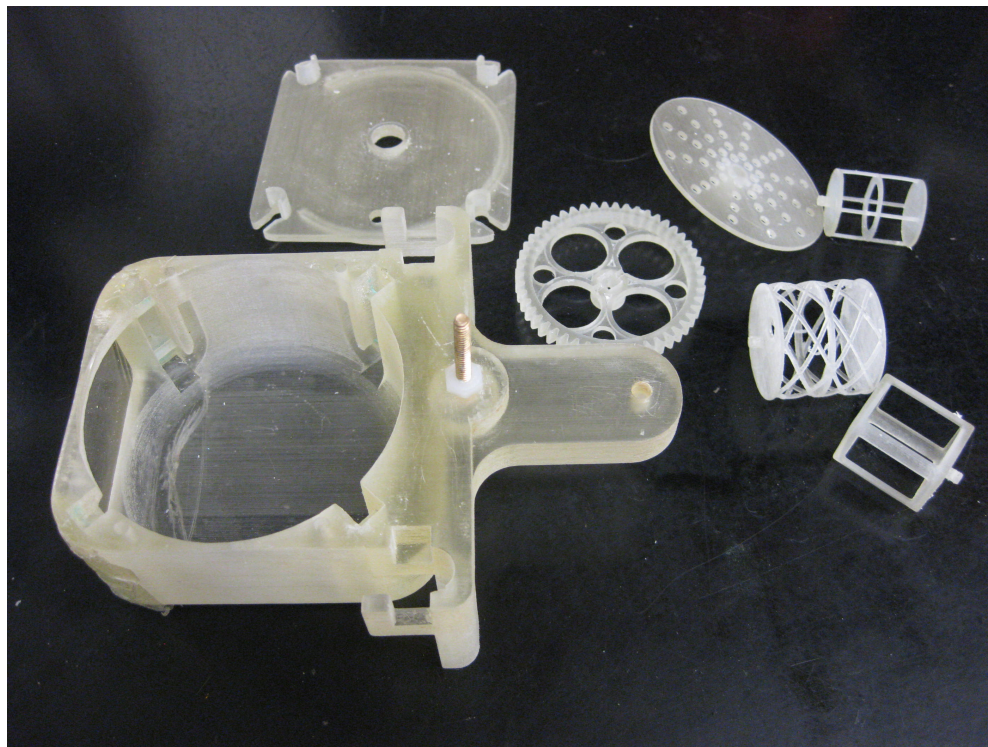
# The experiment

Joint work with **R. Witte** and **P. Ingram**, Medical Imaging Department, UA

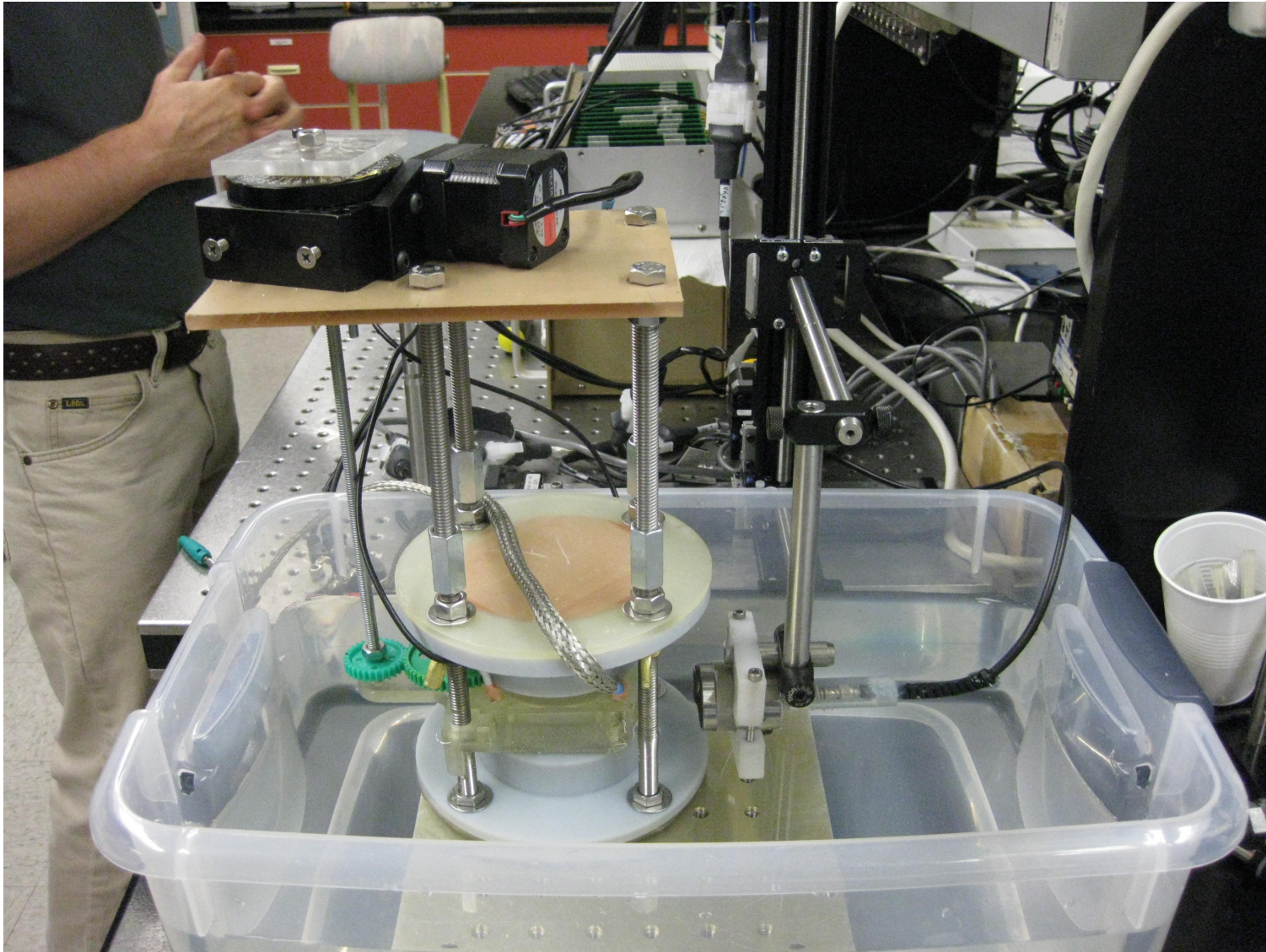
Supported by a BIO5 fellowship, but no money for hardware :(

Goal: build the first MAET scanner, get first MAET images

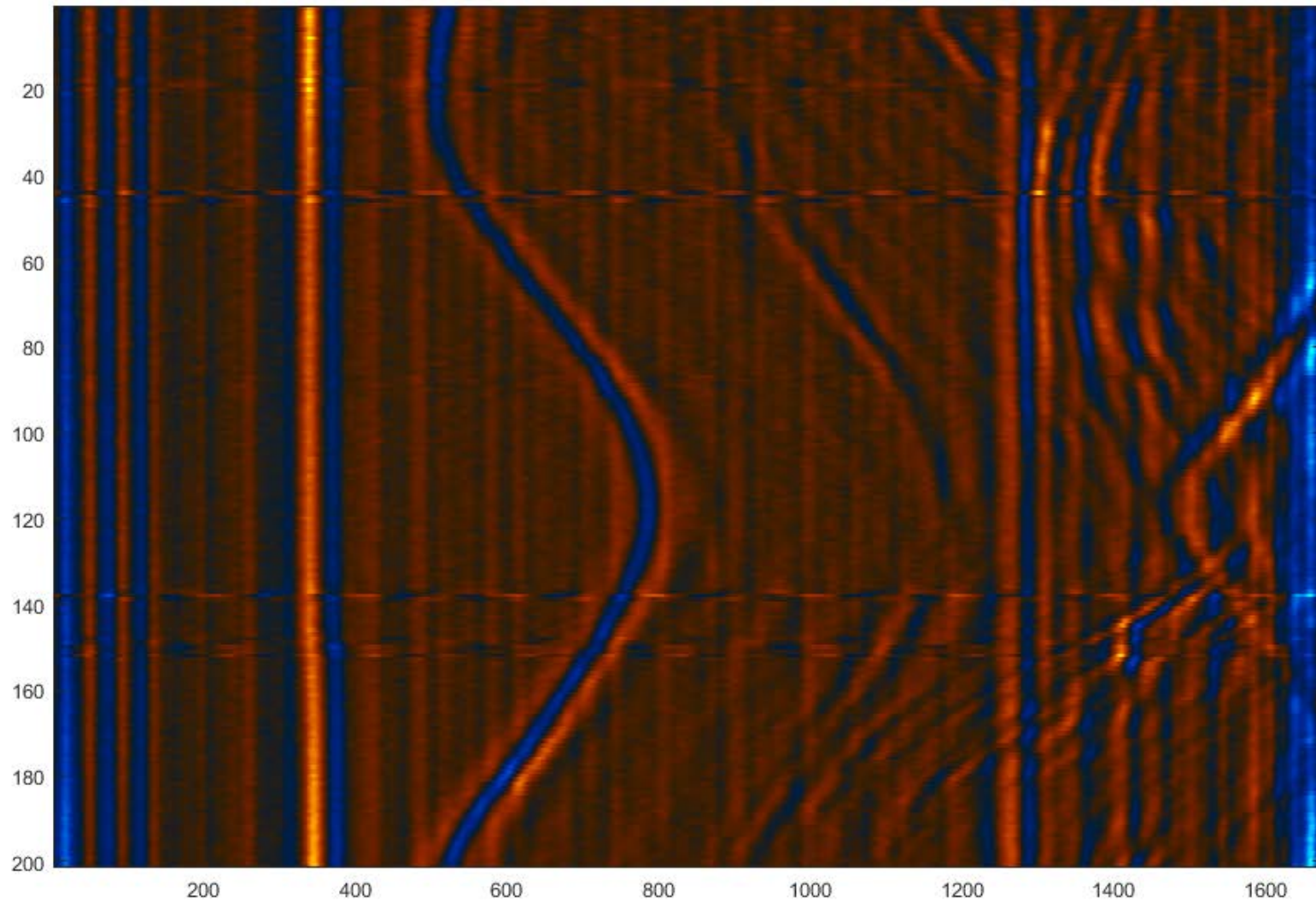
Parts were designed in SolidWorks and 3D-printed



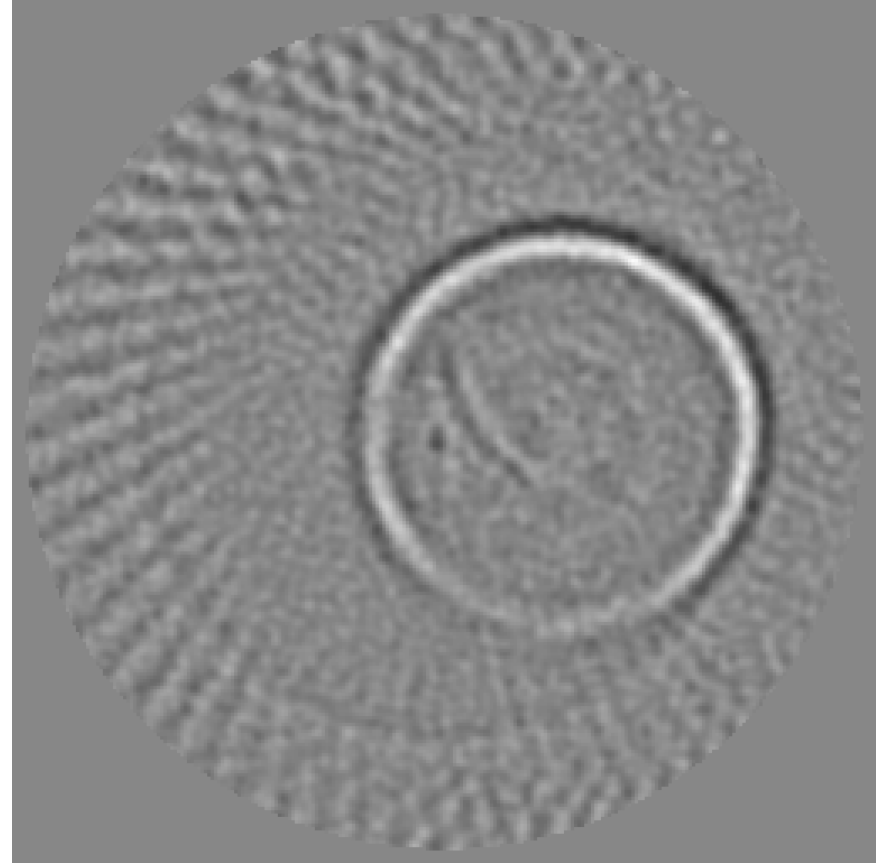
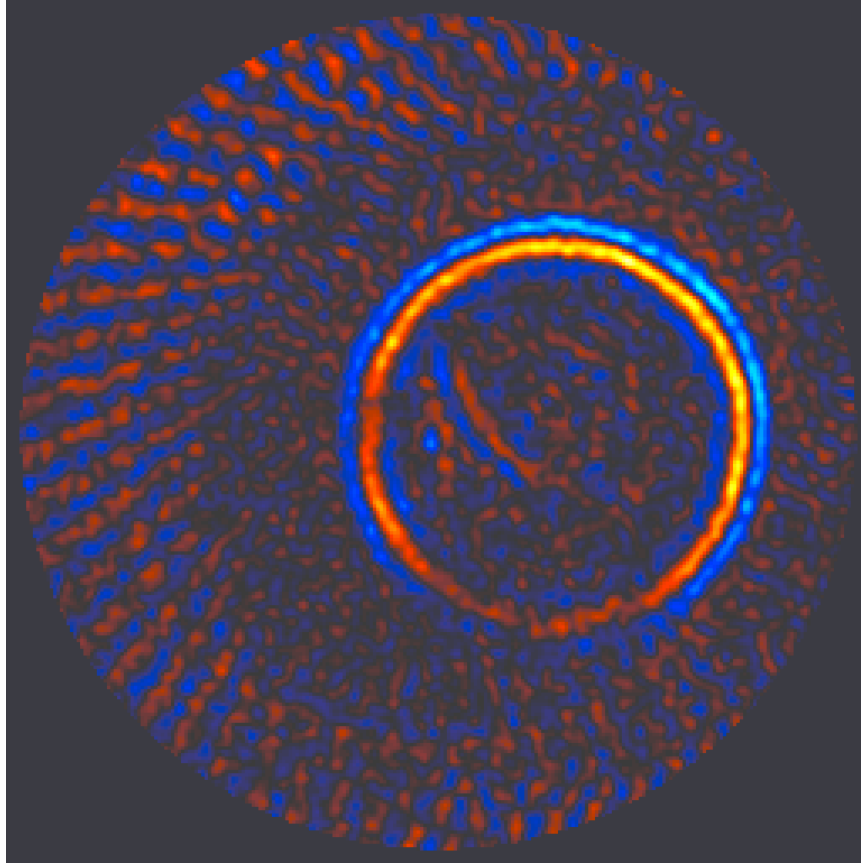
# Fully assembled, in a tank, with a transducer



# How does the signal look?

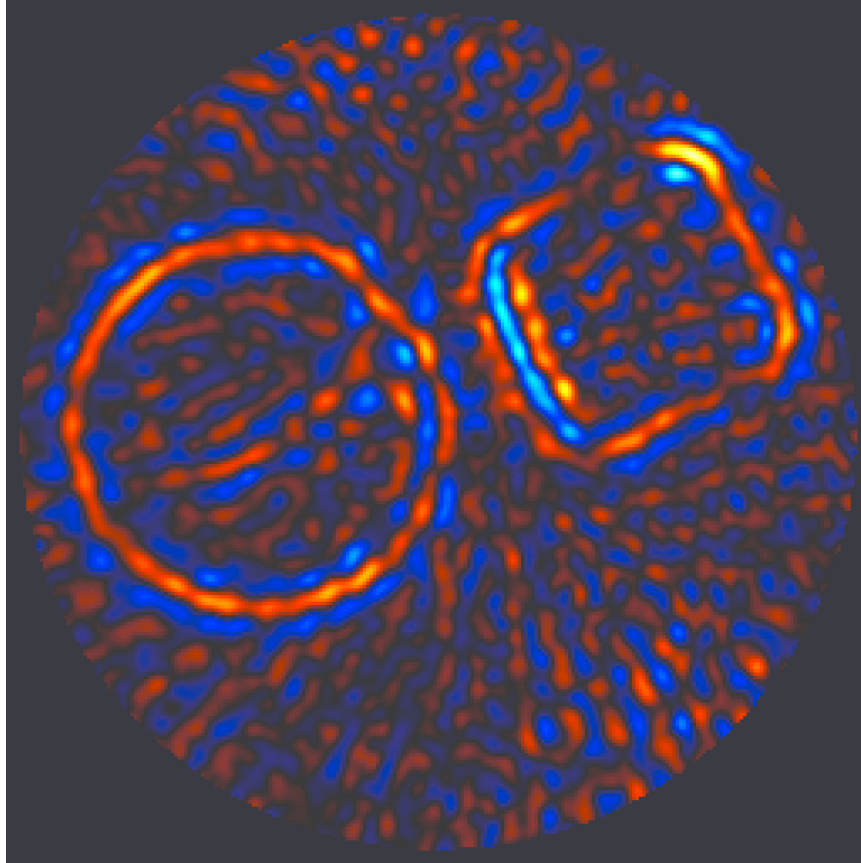


# First reconstruction: round non-conducting phantom

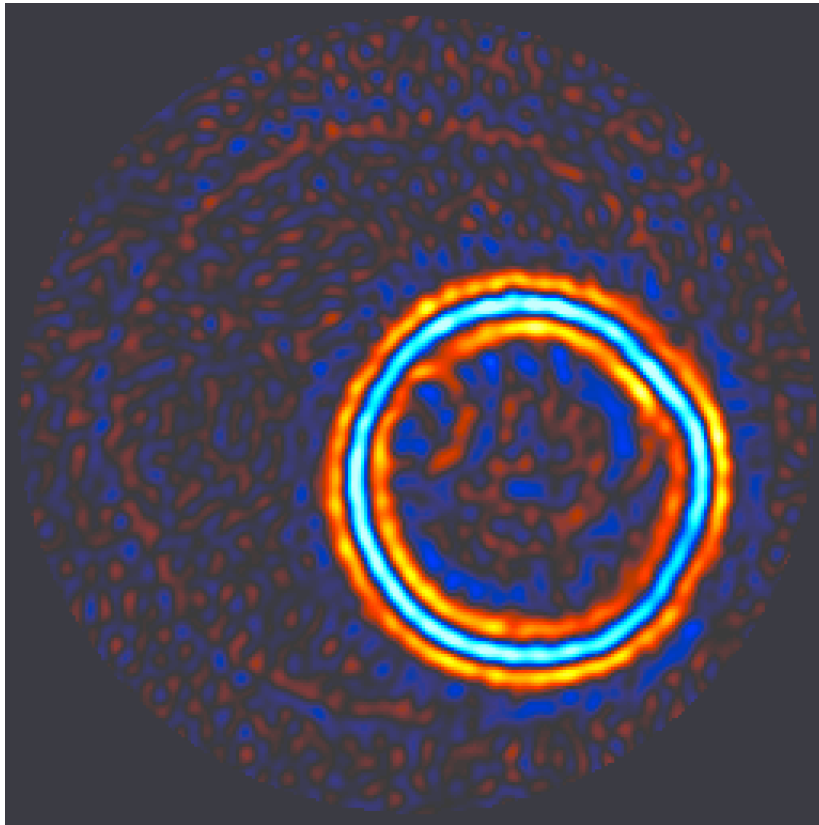
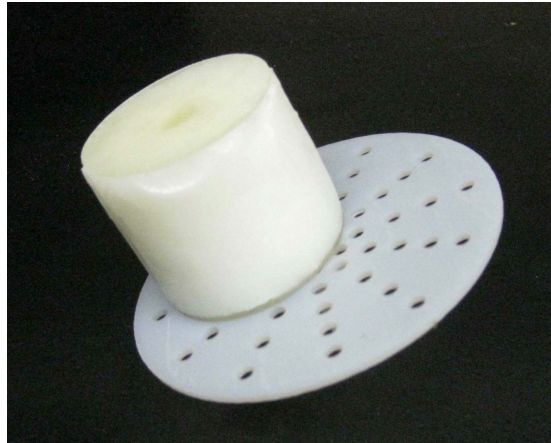




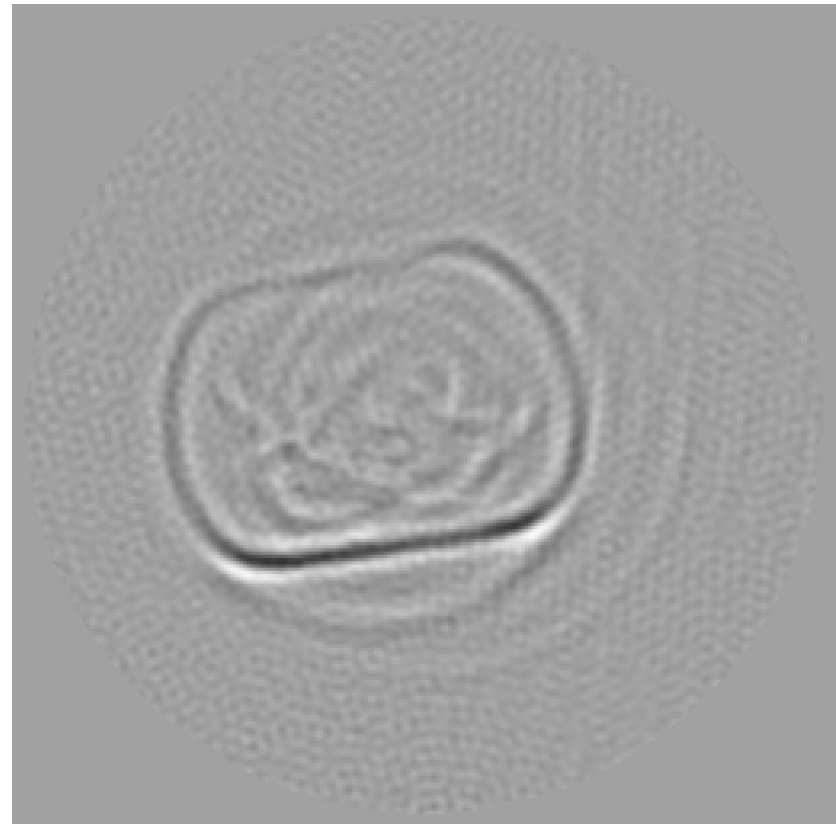
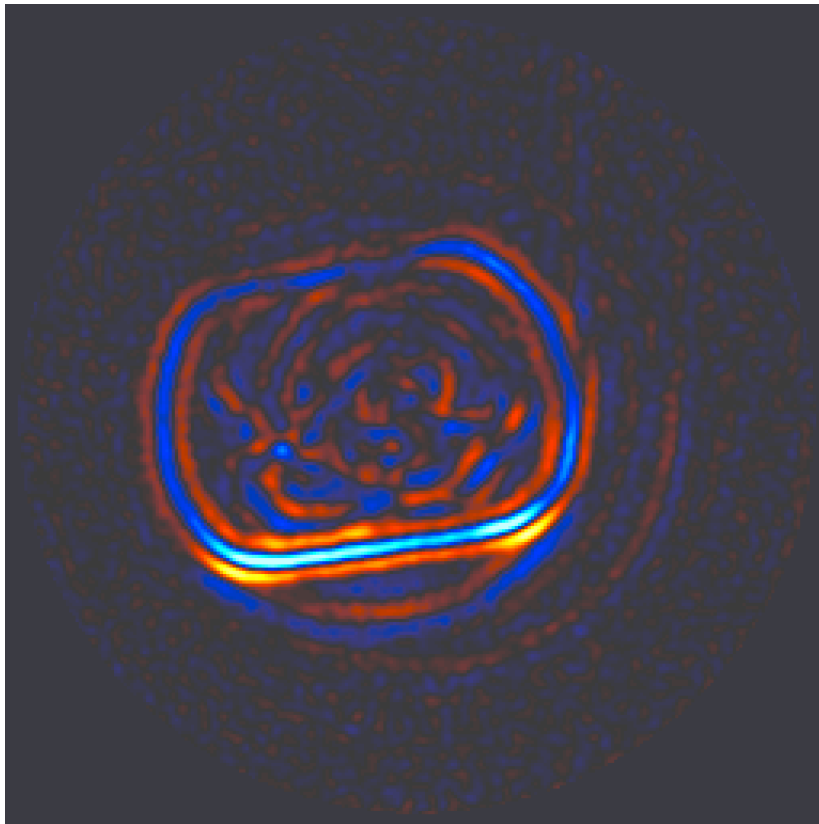
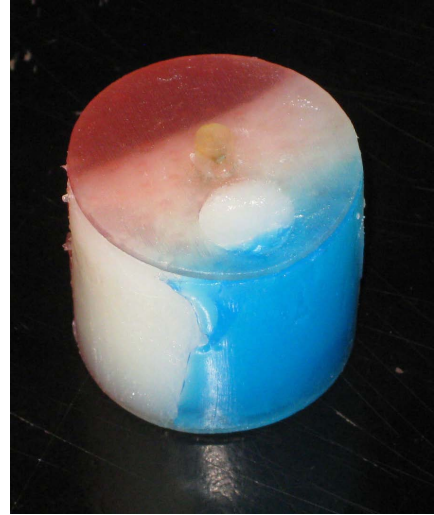
# Round and square non-conducting phantoms



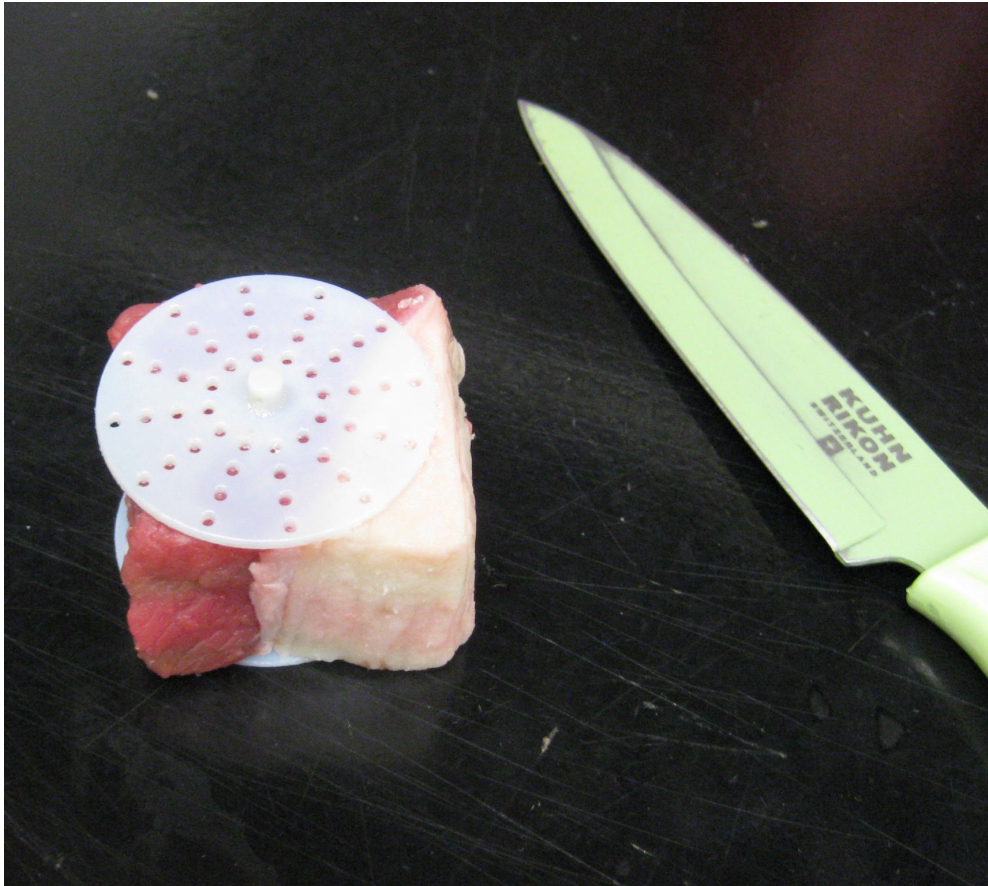
# Round lard column, 30mm in diameter



# Layered gel-lard-gel phantom



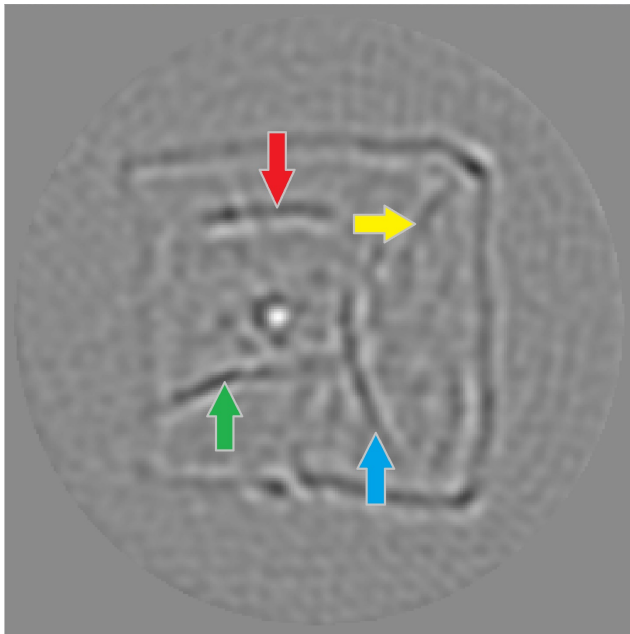
# A beef sample



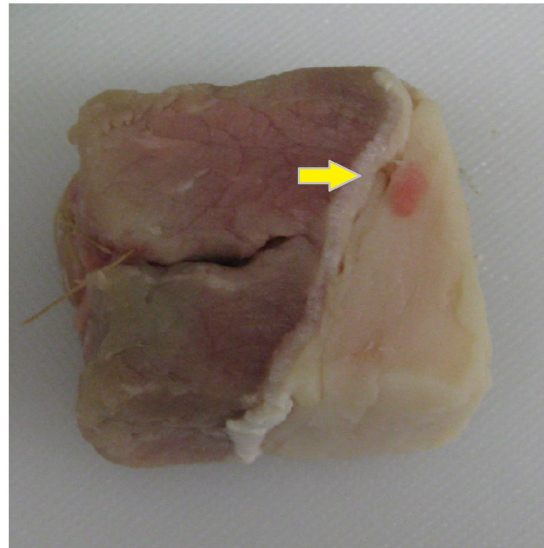
This is **the first** truly tomographic MAET image of a biological tissue.

# Are the details in the image real?

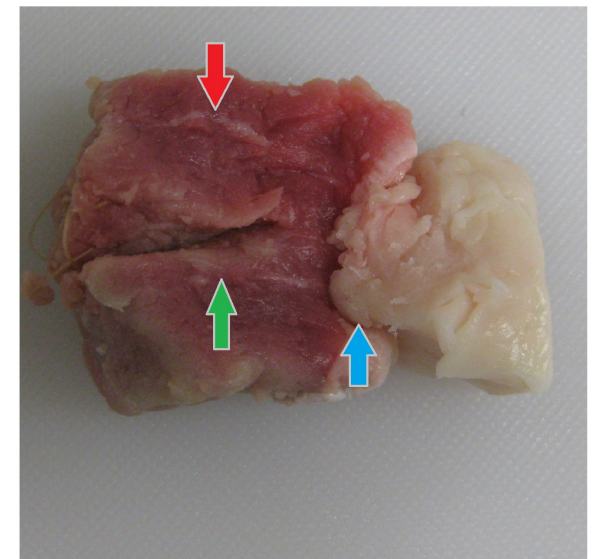
Image



Sample



Sample cut  
in half



# What's next?

Photoacoustic generation of ultrasound waves?

MAET in a bore of an MRI scanner?

**The end**